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Polarization Observables in Pion Photoproduction:
 A Coupled Channels Model Predictions

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Δ -isobar plays a significant role in medium energy π -nucleus reactions. Coupled channels model¹⁾ (CCM), which enlarges the Hilbert space in conventional nuclear many-body theory to include explicitly the Δ degrees of freedom, has been constructed to deal with this fact. Such a model has now been applied to NN and πd scattering, pion absorptions by deuteron, pion production in NN collision with considerable success. We have recently extended²⁾ the CCM to photopion processes in the hope that it may serve as a unified theoretical framework for π -nucleus and photopion production reactions.

For the simpler cases of πN scattering and $\gamma N \rightarrow \pi N$, the Hamiltonian of such an unified model can be written as

$$H = H_0 + H_{\pi N \Delta} + H_{\gamma N \Delta} + \bar{V}_{\pi N} + V_{\gamma \pi}^{(B)}, \quad (1)$$

where H_0 is the sum of relativistic free-energy operators for N, Δ , π and γ . $\bar{V}_{\pi N}$ is the πN interaction except in P_{33} channel. $V_{\gamma \pi}^{(B)}$ is the transition potential operator whose matrix elements will give rise to a transition amplitude for $\gamma N \rightarrow \pi N$ as given by standard effective Lagrangian approach³⁾ which includes the Born terms in pseudovector coupling and contributions from t-channel (ω, ρ) vector-meson exchanges. $H_{\pi N \Delta}$ and $H_{\gamma N \Delta}$ are the $\pi N \leftrightarrow \Delta$ and $\gamma N \leftrightarrow \Delta$ vertex interactions, respectively, in P_{33} channel. The transition matrix for $\gamma N \rightarrow \pi N$, to first order in e , is then given by

$$t_{\gamma \pi} = t_{\pi N} G_0 V_{\gamma \pi} + V_{\gamma \pi}, \quad (2)$$

with

$$t_{\pi N} = t_{\pi N} G_0 V_{\pi N} + V_{\pi N} \quad (3)$$

$$V_{\pi N} = \bar{V}_{\pi N} + H_{\pi N \Delta} (E - E_{\Delta})^{-1} H_{\pi N \Delta} \quad (4)$$

$$V_{\gamma \pi} = V_{\gamma \pi}^{(B)} + H_{\pi N \Delta} (E - E_{\Delta})^{-1} H_{\gamma N \Delta}. \quad (5)$$

$G_0 = (E - H_0)^{-1}$ is the propagator and E_{Δ} denotes the free energy of Δ . We emphasize

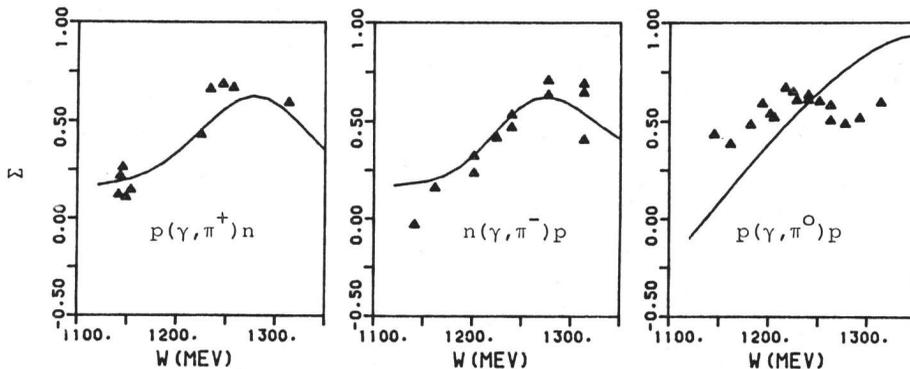


Fig. 1. Our model predictions for polarized photon asymmetry Σ at $\theta=90^\circ$.

that the $\gamma\pi$ amplitude thus obtained above satisfies the Watson's theorem because the unitarity has been built into the model by treating microscopically the pion multiple scatterings. This is a feature which distinguished our model from others constructed before⁴⁾.

With proper choices of $\bar{V}_{\pi N}$ and coupling constants in $H_{\pi N\Delta}$ and $H_{\gamma N\Delta}$, good fits to πN phase shifts and $\gamma\pi$ multipole amplitudes in all channels have been obtained²⁾. It is known that the polarization data provide a severe testing ground for the complicated spin and isospin structure of the photopion production amplitudes, and we present here the predictions of our model for these quantities. Fig.1 gives the results for polarized photon asymmetry Σ at $\theta=90^\circ$ for reactions $p(\gamma,\pi^+)n$, $n(\gamma,\pi^-)p$ and $p(\gamma,\pi^0)p$, respectively. The results for recoil nuclei polarization P and polarized target asymmetry T , both at $\theta=90^\circ$, are shown in Fig.2 for $p(\gamma,\pi^+)n$. The results shown here are obtained without the (ω,ρ) contributions in $V_{\pi}^{(B)}$. Nevertheless, the agreement with experiments are impressive except in the case of Σ for $p(\gamma,\pi^0)p$. It is expected that the agreement will further be improved after we finish the full calculation with the inclusion of (ω,ρ) contributions.

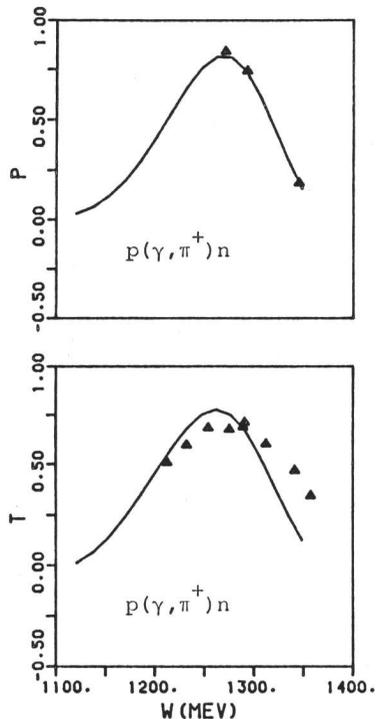


Fig. 2. Our model predictions for polarized target asymmetry T and recoil nuclei polarization P at $\theta=90^\circ$.

References

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