

### 7.13 Hyperfine Interactions of Spin Polarized $\beta$ -Emitter $^{12}\text{N}$ in Fe Crystal

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Hyperfine fields detected by dilute  $^{12}\text{N}$  impurity at the interstitial sites of ferromagnetic Fe were studied as well as its lattice distortion due to the impurity. The technique employed here was essentially the same with that used in the previous works<sup>1,2)</sup>, i.e. the NMR detection in which the asymmetric  $\beta$  decay from the polarized  $^{12}\text{N}$  ( $I^\pi=1$ ,  $T_{1/2}=11\text{ms}$ ) nuclei was monitored<sup>3)</sup>. Nuclear spin polarized  $^{12}\text{N}$  was obtained at recoil angle  $\phi = 15\sim 25$  degrees in  $^{10}\text{B}(^3\text{He},n)^{12}\text{N}$  reaction initiated with  $^3\text{He}$  ions of 3.0 MeV. The polarization detected in the  $\beta$ -decay asymmetry was  $P = 0.25$ . The angular distribution of the  $\beta$ -ray is given by  $W(\theta) = 1 + P\cos\theta$  where  $\theta$  is the polar angle of  $\beta$ -ray momentum relative to the polarization direction. A strong magnetic field  $H_0$  was applied parallel to the polarization for the purpose of the NMR detection, which field was also very effective to maintain the polarization during the flight of the  $^{12}\text{N}$  ions in vacuum before they reached an implantation material. The ions ejected from the target were embedded (implanted) in Fe crystal by use of the kinetic energy obtained in the nuclear reaction. Because of the heterogeneous energy spread due to the target thickness of about  $100 \mu\text{g}/\text{cm}^2$  and the spread in the recoil angle, the distribution of  $^{12}\text{N}$  ions in the sample was almost even throughout the region from the surface to the maximum recoil range. Since the estimated density of  $^{12}\text{N}$  was very dilute,  $\sim 10^{10}/\text{cc}$ , only the interaction of the present  $^{12}\text{N}$  impurity with host atoms were studied, in other words, the interaction with other  $^{12}\text{N}$  impurities themselves was negligible.

A typical NMR spectrum, the polarization change, as a function of rf frequency at  $H_0 = 7 \text{ kOe}$  and  $T = 120 \text{ K}$ , is shown in Fig. 1. The crystal axis  $\langle 011 \rangle$  was set parallel to  $H_0$ . The observed field consisted of five components,

$$\vec{B}_{\text{obs}} = \vec{B}_{\text{ext}} + \vec{B}_L + \vec{B}_{\text{DM}} + \vec{B}_{\text{hf}} + \vec{B}_{\text{dip}} + \vec{\omega}_Q/\gamma,$$

where  $B_i$  in the right hand side are external magnetic field, Lorentz field, demagnetization field, hyperfine field, and dipolar field.  $\omega_Q$  is the angular frequency due to the nuclear quadrupole interaction. Since  $B_L$  and  $B_{\text{DM}}$  are estimated from the known magnetization of Fe at temperature  $T$ , i.e.  $M_s(T)$ , and also  $\omega_Q$  is measured in the experiment ( $B_{\text{hf}} + B_{\text{dip}}$ ) is easily determined from the NMR spectra.

From the NMR spectra as a function of crystal orientation relative to  $H_0$ , two locations were determined. For the resonance line at 1.9 MHz,  $B_{\text{dip}}$  due to the possible anisotropic distribution of surrounding Fe atoms around an interstitial  $^{12}\text{N}$  is deduced to be  $B_{\text{dip}} \sim 0$ . Also this was well explained by two lines split due to the quadrupole interaction of  $\omega_Q/2\pi = 0.75eqQ/h = 400 \text{ kHz}$ . The lines are from  $^{12}\text{N}$  in an interstitial site. Experimental hyperfine field due to the contact interaction of  $^{12}\text{N}$  nucleus with conduction electrons of S-symmetry,  $B_{\text{hf}} = -(9.41 \pm 0.11) \text{ kOe}$  was obtained. For the resonance at 3.7 MHz, finite amount of dipolar field  $B_{\text{dip}} \sim \pm 2 \text{ kOe}$  was observed, and the hyperfine field of  $-35 \text{ kOe} < B_{\text{hf}} < +20 \text{ kOe}$  was determined. In spite of the appreciable amount of  $B_{\text{dip}}$ , we could not extract it clearly neither determine the location of  $^{12}\text{N}$  from this. The sign of  $B_{\text{hf}}$  was not also determined because of the unresolved  $B_{\text{dip}}$  and  $\omega_Q$ . Although, we can conclude that the locations are interstitial, exact sites were not determined. However, those observed  $B_{\text{hf}}$  values are extraordinary small compared with theoretical one based on ab initio band calculation<sup>4)</sup> for  $^{12}\text{N}$  ions in tetrahedral and octahedral sites by J. Kanamori et al., i.e.  $B_{\text{dip}} \sim -70 \text{ kOe}$ . If, in this powerful theory, the lattice renormalization (expansion) of about  $\Delta a/a = 50\%$  for, at least, the nearest surroundings, which was observed for  $^{12}\text{N}$  in V crystal<sup>3)</sup> taken into account, a better agreement in the experimental and theoretical values will be obtained. It is also pointed out in this

hyperfine interaction studies of  $^{12}\text{N}$  in Fe that they are implanted in two independent interstitial sites. This shows a strange deviation from the systematics on  $^{12}\text{N}$  ions in other crystals that they are located in an interstitial site.

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#### References

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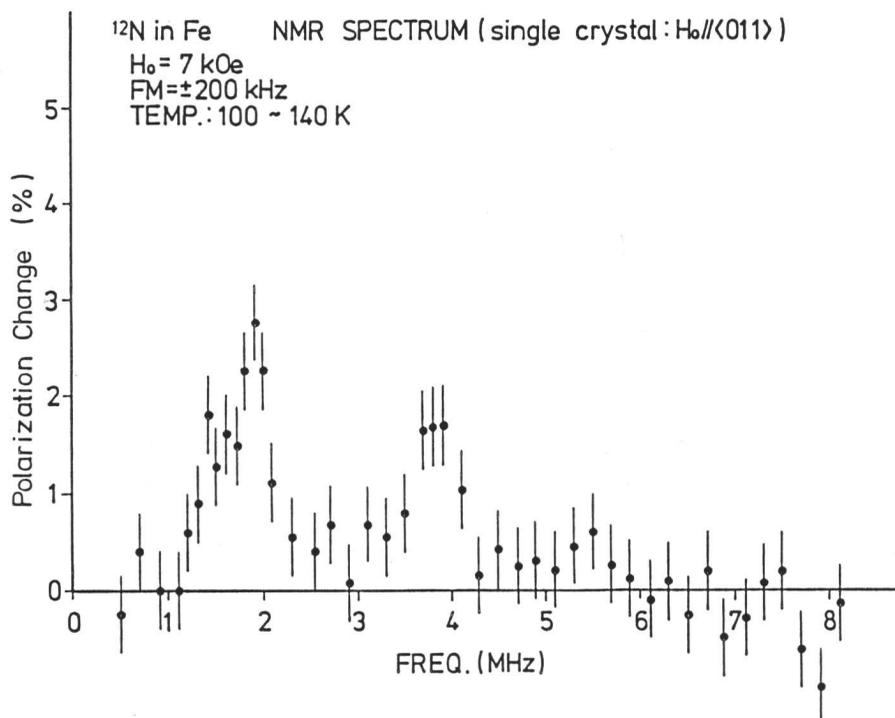


Fig. 1. NMR spectrum of  $^{12}\text{N}$  in ferromagnetic Fe.  
 The crystal axis  $\langle 011 \rangle$  of Fe was parallel to  $H_0$  of 7.0 kOe at  $T = 120 \text{ K}$ .